

that the "resonance" occurs when one of the higher harmonics $\cos nx$ is associated with the same linearized wave speed as the (fundamental) $\cos x$. Now, from our earlier linear results or from (41) with $\epsilon = 0$, we know that a wave of length $2\pi L$ moves with speed C given by

$$C^2 = gL \left(1 + \frac{T}{\rho g L^2} \right). \quad (45)$$

A harmonic of wavelength $1/n$ times as large moves at the same speed if

$$\frac{gL}{n} \left(1 + \frac{n^2 T}{\rho g L^2} \right) = gL \left(1 + \frac{T}{\rho g L^2} \right) \quad \text{or} \quad \frac{T}{\rho g L^2} = \frac{1}{n}, \quad n = 2, 3, \dots$$

This condition, $B = n^{-1}$, is indeed the one that holds when our original perturbation calculations fail.

Consider further the case $B = \frac{1}{2}$; i.e., $L = \bar{L}$. According to linear theory, the periodic wave

$$a \cos x + b \cos (2x + \theta) \quad (46)$$

will travel without change of shape, whatever the values of the amplitudes a and b and the phase factor θ . The reason is of course that each of the two constituent waves travels at the same speed (about 24 cm/s), in spite of the fact that one is twice as long as the other. This suggests that (46) is the proper solution to assume at lowest order when $B = \frac{1}{2}$. Such an assumption was made in a paper by J. R. Wilton in 1915, and later, in a slightly modified fashion, in an article by W. J. Pierson and P. Fife.* The solution process can be carried forward without contradiction. The results for the surface deflection as given by Pierson and Fife is (in our notation)

$$\epsilon \zeta = \epsilon [\cos x + (\pm \frac{1}{2} + \epsilon b_2) \cos 2x] \mp \frac{3}{2} \epsilon^2 \cos 3x - \frac{1}{2} \epsilon^2 \cos 4x + \dots \quad (47)$$

There are two different possible solutions when $B = \frac{1}{2}$. The solution corresponding to the upper (lower) signs is of the gravity (capillary) type, with "dimples" up (down). The coefficient b_2 is not determined at this order. As pointed out by L. F. McGoldrick, however,† constants like b_2 are ultimately determined by initial conditions, so they can be chosen at will to improve the agreement with any given experimental data.

The dimensional wave speed corresponding to (47) is given by

$$C = \frac{3}{2} gL \left(1 \pm \frac{\epsilon}{4} + \dots \right). \quad (48)$$

Noteworthy is the presence of a first order correction term.

* "Some Nonlinear Properties of Long-Created Periodic Waves with Lengths Near 2.44 Centimeters," *J. Geophys. Res.* 66, 163-79 (1961).

† "On Wilton's Ripples: A Special Case of Resonant Interactions," *J. Fluid Mech.* 42, 193-200 (1970).

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Michael Hoelscher



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